Radiative signatures of magnetic reconnection in blazar jets

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Main ingredients

1. Dissipation

- Close vs. far from SMBH
- Magnetic- vs. kineticdominated jet
- Structured vs. uniform jet
- · Continuous vs. localized





Leon-Tavares et al. 2014





2. Radiation

- Radiating particle species
- Radiation processes
- Properties of accelerated particle distributions



Magnetic reconnection



- Dissipation of magnetic energy
- Transformation into heat, bulk plasma kinetic energy and non-thermal particle energy
- Magnetic kink instability at jet interior
 - Striped wind structure of jet

(e.g. Romanova & Lovelace 1992; Eichler 1993; Begelman 1998; Giannios & Spruit 2006; McKinney & Uzdensky 2012, Yuan et al. 2016; Blandford+2017 (review); Alves et al. 2018; Giannios & Uzdensky 2019)

A radiative model of magnetic reconnection



Reconnection dynamics, energetics, and particle energy spectrum

Appearance of the reconnection layer

2D PIC simulations in electron-positron plasma with σ =10



- Inflow into the layer is non-relativistic, $|v_{in}|/c \sim 0.10$ (Lyutikov & Uzdensky 03, Lyubarsky 05)
- Outflow from the X-points is relativistic, reaching the Alfvén speed, $v_A/c = \sqrt{\sigma/(\sigma+1)}$

• Current sheet breaks into a series of magnetic islands or plasmoids (e.g. Loureiro+07, Bhattacharjee+09, Uzdensky+10, Huang & Bhattacharjee 12).

Particle energy spectrum

2D σ =10 electron-positron

3D σ =10 electron-positron



- Higher σ produces harder (p<2) particle spectra. (Sironi & Spitkovsky 2014; Guo+2014, 2015; Werner+2016)
- The maximum energy increases with time, if it is not inhibited by the boundaries. (Sironi & Spitkovsky 2014; Petropoulou & Sironi 2018)

Energy dissipation efficiency



Efficiency:

$$f_{\rm rec} \equiv \frac{\sum_i \int_{V_i} U_{\rm e} dV_i}{\sum_i \int_{V_i} (e + \rho c^2 + U_{\rm B}) dV_i}$$

- Reconnection transfers ~ 50% of the flow energy (electron-positron plasmas) or ~ 25% (electron-proton) to radiating particles.
- Efficiency decreases with increasing guide field.

Sironi, Petropoulou, Giannios 2015 (see also Werner+2017)

Energy equipartition



Weighted average energy density ratio:

$$\left\langle \frac{U_{\rm e}}{U_{\rm e} + U_{\rm B}} \right\rangle \equiv \frac{\sum_{i} \int_{V_{i}} U_{\rm e} \frac{U_{\rm e}}{U_{\rm e} + U_{\rm B}} \mathrm{d}V_{i}}{\sum_{i} \int_{V_{i}} U_{\rm e} \mathrm{d}V_{i}}$$

- Rough energy equipartition in plasmoids: even if the upstream plasma is strongly magnetized, the "effective" σ ~1 within plasmoids.
- $U_{e} >> U_{B}$ is not expected.
- $U_{\rm B} >> U_{\rm e}$ for strong guide fields.

Plasmoid growth & acceleration



- The plasmoid size grows with constant rate $\beta_g \sim 0.1$ (in its rest frame).
- The growth rate depends weakly on the magnetization σ .

Sironi, Giannios, Petropoulou 2016



• "Universal" acceleration profile due to magnetic forces:

$$\Gamma \frac{v_{\text{out}}}{c} \simeq \sqrt{\sigma} \tanh\left(\frac{\beta_{\text{acc}}}{\sqrt{\sigma}} \frac{x - x_0}{w}\right)$$

- Small plasmoids are relativistic.
- Biggest plasmoids move with non-relativistic speeds.

Plasmoid statistics

Distribution of sizes



Distribution of 4-velocities



Petropoulou, Christie, Sironi, Giannios 2018

- Slope of power-law size distribution depends on plasmoid growth & acceleration rates.
- 4-velocity distribution depends on σ .

Radiative signatures from plasmoids

The radiative model





- Use 2D PIC of magnetic reconnection to benchmark:
 - i. the particle energy spectrum (slope, minimum, maximum energies),
 - ii. the growth rate of emitting regions,
 - iii. the ratio of electron-to-magnetic energy density, and
 - iv. the Doppler factor of emitting regions.
- Vary free parameters to produce generic blazar spectra:
 - i. magnetic field strength and magnetization σ ,
 - ii. length & orientation of the layer,
 - iii. pair multiplicity,
 - iv. jet's bulk Lorentz factor, and
 - v. external radiation fields.

Emission from a single plasmoid

Each plasmoid produces a flare of characteristic duration and flux

 $\Delta t_{1/2} \approx \frac{\mathbf{w}_p}{\beta_g c \, \boldsymbol{\delta}_p}$

 $L_{pk} \approx \frac{f_{rec} L_j}{8 R^2 c \beta_j \Gamma_j^2} \beta_g c w_p^2 \delta_p^4$

Small & fast → Minute-scale flare

Large & slow \rightarrow Hour-scale flare



Emission from the whole layer

Each reconnection event produces a chain of plasmoids with power-law size distribution



- Fast flares on top of slowly evolving envelope (Giannios 2013).
- Physical model for multi-timescale variability in jets.

Flux-flux correlations



BL Lac objects and FSRQs



(BL Lac)

- Different plasma magnetization $\sigma \rightarrow$ power-law index of electron distribution
- Different pair multiplicities $\rightarrow \gamma_{min}$ for soft PL (γ_{max} for hard PL)

Not yet demonstrated with PIC

Luminosity of broad line region \rightarrow strength of external Compton



Effects of pair multiplicity

First 2D PIC simulations in electron-positron-proton plasmas with σ =1, 3,10, Θ_{ρ} =0.1, 1, 10 and pair multiplicity κ =1-199



Empirical scaling of mean lepton Lorentz factor on physical parameters:

$$\langle \gamma_e - 1 \rangle \approx \sqrt{\sigma} (1 + 4\Theta_e) \left(1 + \frac{\sigma_{e,h}}{30} \right)$$

magnetization of pairs defined with enthalpy



- Higher $\kappa \rightarrow$ softer power laws
- Higher $\kappa \rightarrow$ lower mean lepton energy
 - Larger $\sigma \rightarrow$ harder power laws

Inter-plasmoid Compton scattering



Christie et al. 2019, MNRAS submitted (arXiv:1908.02764)

- Large plasmoids can provide extra seed photons for Compton scattering.
- Inter-plasmoid Compton scattering (IPCS) is important for a large fraction of plasmoids in layer.
- IPCS in BL Lacs can explain A_c~ 1-5 without requiring particle-dominated emitting regions!

Fake Fermi light curves from reconnection

Meyer, Petropoulou, Christie, in prep.

Our method:

- i. Select a real FSRQ and a random time window.
- ii. Conduct standard *Fermi* analysis of optimizing source parameters in a given ROI, and build a light curve (with fixed temporal binning).
- iii. For each time bin, replace the central source by the prediction of our "vanilla" blazar model, add a baseline flux, and simulate the modified ROI.
- iv. Conduct a standard *Fermi* analysis on the simulated ROI and reconstruct the source light curve.

Our goal:

i. Test which (if any) features of the model light curves are retained in a *Fermi*-LAT observation.



Summary

Magnetic reconnection:

- is an efficient dissipation process.
- produces extended non-thermal particle energy spectra.
- leads to the generation of many plasmoids, aka emitting regions.
- can produce multi-wavelength non-thermal emission.
- can produce variable emission on multiple timescales.
- offers a physical framework for jet emission.

Future directions

- Systematic study of flare statistics: flux distributions, power spectral density, flux-flux correlations.
- Neutrino emission from reconnection layers.
- Connection to large-scale jet properties: Where? How? Guide field?



THANK YOU!

Back-up slides

Plasmoid compression: a slow acceleration process

2D σ =10 electron-positron



Effects of layer's orientation



Christie et al. 2019

Energy equipartition in pair-proton reconnection



Standard blazar SED modeling

- Both models describe equally well steady-state & timeresolved photon spectra.
- The models imply different physical conditions for the radiation zone.



3C273

Leptonic

 10^{1}

 10^{13}

Modeling studies: Krawczynski et al. 2002; Ghisellini et al. 2010; Boettcher et al. 2013; Potter & Cotter 2013; Dimitrakoudis, MP, Mastichiadis 2014; MP 2014; MP, Dimitrikoudis et al. 2015; Cerruti et al. 2015; Diltz, Boettcher & Fossati 2015; Gao et al. 2017; MP, Nalewajko et al. 2017; Cerruti et al. 2017 +++

